

The Rendering Equation

Kajiya '86

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Since Derek made slides. . .

$$I(x, x') = g(x, x') \left[\epsilon(x, x') + \int_S \rho(x, x', x'') I(x', x'') dx'' \right]$$

- I : intensity of light from x to x'
- g : geometry
- ϵ : emission
- ρ : intensity of light from x'' to x from x'

- Time-averaged transport intensity
 - No phasees
 - No diffraction
- Homogeneous refractive index of base medium
 - Could let g handle this (single scattering)
 - Could solve this equation volumetrically (Eikonal)
 - Could solve this equation with sparse paths (path integration)

- How should we set I (that didn't exist before this paper)?
- Isn't it more natural to talk about angles around a point?
- Use stoichiometry to connect with standard radiometric intensity!

$$dE = I(x, x') dt dx dx'$$

- $dE \sim \text{joule}$
- $dx, dx' \sim m^2$
- $dt \sim \text{sec}$
- $\Rightarrow I(x, x') \sim \frac{\text{joule}}{m^4 \text{sec}}$

$$g(x, x') \sim \frac{\textit{Visible}(x, x')}{r^2}$$

- Unitless!

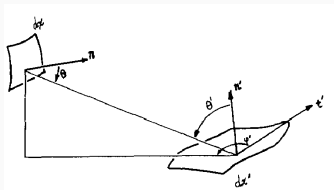
$$dE = \frac{1}{r^2} \epsilon(x, x') dt dx dx'$$

- $dE \sim \text{joule}$
- $dx, dx' \sim m^2$
- $dt \sim \text{sec}$
- $\Rightarrow I(x, x') \sim \frac{\text{joule}}{m^2 \text{sec}}$

$$dE = \frac{1}{r^2} \rho(x, x', x'') I(x', x'') dt dx dx' dx''$$

- $dE \sim \text{joule}$
- $dx, dx', dx'' \sim m^2$
- $dt \sim \text{sec}$
- $I(x, x') \sim \frac{\text{joule}}{m^4 \text{sec}}$
- \Rightarrow unitless!

- $i(\theta, \phi)$
 - per time (dt)
 - per unit projected area (dx'_p)
 - per unit of solid angle ($d\omega$)



$$dE = i(\theta', \phi') d\omega dx'_p dt$$

- After defining maps between points and angles...

$$r = \|x - x'\|$$

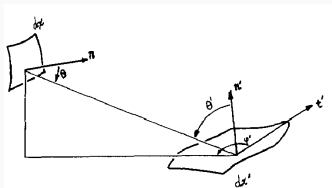
$$dx'_p = dx' \cos \theta$$

$$\cos \theta = \frac{1}{r} \langle n, x - x' \rangle$$

$$\cos \theta' = \frac{1}{r} \langle n', x - x' \rangle$$

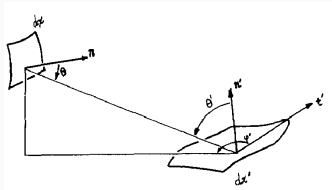
$$\cos \phi' = \frac{1}{r} \langle t', x - x' \rangle$$

$$dE = i(\theta', \phi') \frac{\cos \theta \cos \theta'}{r^2} dt dx dx'$$



- After defining maps between points and angles. . .

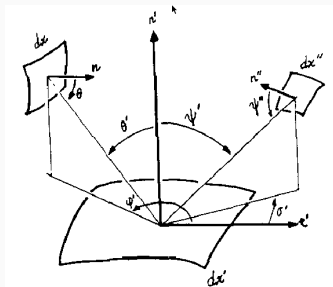
$$I(x, x') dt dx dx' = dE = i(\theta', \phi') \frac{\cos \theta \cos \theta'}{r^2} dt dx dx'$$
$$I(x, x') = i(\theta', \phi') \frac{\cos \theta \cos \theta'}{r^2}$$



- After some more boring algebraic manipulation with several variables. . .

$$\epsilon(x, x') = \epsilon(\theta', \phi') \cos \theta \cos \theta'$$

$$\rho(x, x', x'') = \rho(\theta', \phi', \psi', \sigma') \cos \theta \cos \theta'$$



- The basic rendering equation is really elegant and now we have a bunch of cos everywhere, which is important
- Now we can try solving it maybe?

- As the equation is an affine transformation we can rewrite it as

$$I = g\epsilon + gMI$$

- We can solve a steadystate by...

$$(1 - gM)I = g\epsilon$$

- Taylor expansion of the above provides

$$(1 - gM)^{-1}g\epsilon = g \left(\sum_{i=0}^{\infty} [Mg]^i \right) \epsilon$$

- Intuitively,
 - $i = 0$ is light directly emitted
 - $i = 1$ is light bounced once
 - $i = n$ is what?

- If we have a precomputed illumination ϵ_0

$$I = g\epsilon + gM\epsilon_0$$

- this is too easy?
- Well Whitted proposed something slightly more general

$$I = g\epsilon + g \left(\sum_{i=1}^{\infty} [M_0 g]^i \right) \epsilon_0$$

- M_0 is delta functions with a cosine...

- What if there is no angular dependence?

$$\rho(\theta', \phi', \psi', \sigma') = \rho_0$$

- Well radiosity $dB_r(x)$ (a reduced I turns into

$$dB_r(x') = dx' \rho_0 \pi H(x')$$

- for H hemispherical incident energy per time per area
- With emission we get

$$dB_r(x') = \pi[\epsilon_0 + \rho_0 H(x')] dx'$$

- By energy conservation, $\|M\| \leq 1$.
- Consider

$$x + a + \sum_{k=1}^{\infty} M^k a$$

- We can construct a probability of a path of length n , $\omega \in S^n$ to obtain the contribution from σ_1 to x_ℓ by

$$x_\ell = \left(\prod_{i=0}^{\ell} m_{\sigma_{i-1}, \sigma_i} \right) a_{\sigma_i} \frac{1}{p(\omega)}$$

- Apply splitting to say

$$p(\omega) = p(\omega_0) \prod_{i=0}^{\ell} p(\omega_i, \omega_{i+1})$$

- To not oversample they construct a hierarchy with one sample per cell

